

## F. Einstein (1917) : His A and B coefficients

- Einstein tried to obtain Planck's thermal (black-body) radiation formula by considering
  - Interaction between light (EM waves) in cavity and atoms in wall of cavity
  - thermal equilibrium (statistical physics)
- His work thus came as a result of efforts by
  - Boltzmann (~1900)
  - Planck (~1900)
  - Einstein (1905) (photon)
  - Bohr (~1913) (atom)  
(but prior to quantum theory)

## Planck (1900) Formula

$$U(f, T) df = \frac{8\pi f^2}{c^3} \cdot hf \cdot \frac{1}{e^{\frac{hf}{kT}} - 1} df$$

= energy per unit volume inside cavity due to EM waves  
in freq. range  $f$  to  $f+df$  at equilibrium at temperature  $T$

For our purpose, write as

$$U(\omega, T) d\omega = \frac{\omega^2}{\pi^2 c^3} \cdot \hbar\omega \frac{1}{e^{\frac{\hbar\omega}{kT}} - 1} d\omega \quad (33)$$

- Einstein wanted to find a way to derive Planck's formula, i.e.  
an expression for  $U(\omega, T)$

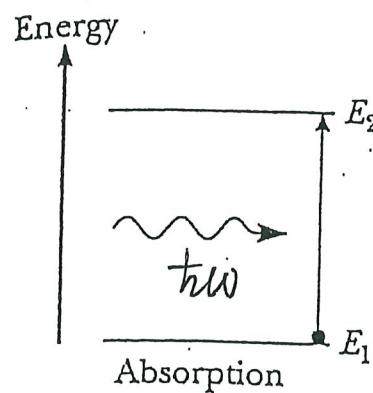
- Bohr (1913)
  - Atom has discrete energies ("stationary states")
  - Spectral lines come from transitions between stationary states
- Einstein (1917)
  - Cavity's wall  $\Rightarrow$  Many Atoms (they absorb/emit; they vibrate)
  - Capturing Bohr's key idea, assumed two-state systems

$$\omega = \frac{E_2 - E_1}{\hbar}$$

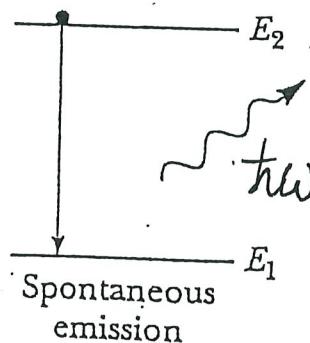
$E_2$  ————— ↑  
↓  $\hbar\omega$  —————  $E_1$

- there could be various values of  $\omega$
- this  $\omega$  will be the " $\omega$ " in  $U(\omega, T)$
- Atoms absorb and emit radiation (exchange energies) from and to the radiation in cavity  $\Rightarrow$  equilibrium can be achieved

To get Planck's formula (especially "-1" in denominator), Einstein needed to invoke 3 processes

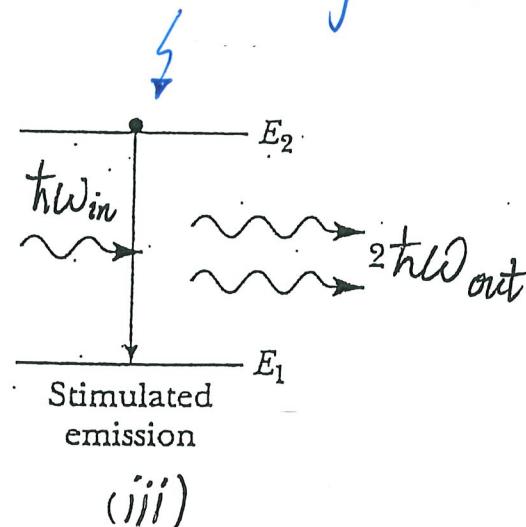


(i)



(ii)

Not known before Einstein's work



(iii)

The processes of (i) absorption and (ii) spontaneous emission. The lifetime of the upper state is  $t_s$ , and the photon is emitted in a random direction.  
 (iii) Stimulated emission. In this process, the emitted photons are in phase with the stimulating photon, and all have the same direction of travel.

- (i) and (ii) were regarded "common sense" processes
- Einstein needed (iii) to obtain Eq.(33)

(i) Stimulated Absorption (level 1 → level 2)

"Stimulated" ⇒ Probability of absorption per atom per unit time  
 $\propto U(\omega, T)$  [this is what "stimulated" meant]

Write: Probability of absorption per atom per unit time  
 $= B_{12} U(\omega, T)$  (34a)

$[B_{12}$  = Einstein's B-coefficient = (simply)  $B]$

Remark: We know from QM that (see Eq.(30))

$$\lambda_{1 \rightarrow 2} = \underbrace{\frac{\pi e^2}{3\epsilon_0 h^2} |r_{21}|^2}_{B_{12} \text{ from QM}} U(\omega) \quad \begin{array}{l} (\text{Einstein didn't know}) \\ \text{this in 1917} \end{array}$$

(ii) Spontaneous Emission (level 2 → level 1)

"Spontaneous" ⇒ Prob. of spontaneous emission per atom per unit time

$$= \underbrace{A_{21}}_{\text{independent of } U(\omega, T)} \quad (\text{Einstein's A-coefficient}) \quad (34b)$$

[independent of  $U(\omega, T)$  [∴ spontaneous]]

Remark: Harder to understand this process within Schrödinger's QM

(iii) Stimulated Emission (level 2 → level 1)

"Stimulated" ⇒ Prob. of stimulated emission per atom per unit time

$$= B_{21} U(\omega, T) \quad [\text{How is it related to } B_{12}?] \quad (34c)$$

[Without this process, could not get Planck's formula]

Remark: We know from QM,  $B_{21} = B_{12}$  ( $\because \lambda_{1 \rightarrow 2} = \lambda_{2 \rightarrow 1}$ )

Einstein didn't know this in 1917

Take stock:

Per atom per unit time

$$\begin{cases} B_{12} \cdot U(\omega, T) & [\text{stimulated absorption}] \\ A_{21} & [\text{spontaneous emission}] \\ B_{21} \cdot U(\omega, T) & [\text{stimulated emission}] \end{cases}$$


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x

Temperature enters via equilibrium statistical physics

$N_1 + N_2$  atoms  $\rightarrow$   $N_2$  atoms in level 2 (higher energy)  
 $\rightarrow$   $N_1$  atoms in level 1 (lower energy)

- Statistical Physics  $\Rightarrow$  At equilibrium  $N_1 > N_2$  (any temperature)

$$\boxed{\frac{N_2}{N_1} = e^{-(E_2 - E_1)/kT} = e^{-\hbar\omega/kT}} \quad (35)$$

- $kT \ll E_2 - E_1$ ,  $N_1 \gg N_2$  ;  $kT \gg E_2 - E_1$ ,  $N_1 \gtrsim N_2$

Aside:

- For a collection of  $N$  entities (independent ones), each one could be at energies  $E_1$  or  $E_2$  (say  $E_2 > E_1$ ), then

$N_2 \propto e^{-E_2/kT}$  and  $N_1 \propto e^{-E_1/kT}$  when system is at thermal equilibrium at temperature  $T$ .

$$\text{So } \frac{N_2}{N_1} = e^{-(E_2 - E_1)/kT} \rightarrow \text{Meaning}$$

Eq. (35)

$N_2 < N_1$   
more entities in lower state

Can't use  $T$  to put more entities in higher state

When  $N_2 > N_1$ , the system is out of equilibrium

- Equilibrium is attained by balancing  $1 \rightarrow 2$  transitions AND  $2 \rightarrow 1$  transitions

$$\begin{array}{c} \text{\# atoms going from} \\ \text{1 to 2 per unit} \\ \text{time by absorption} \\ \left. \vphantom{\frac{\text{# atoms going from}}{\text{1 to 2 per unit}}} \right\} \\ = \end{array} \begin{array}{c} \text{\# atoms going from} \\ \text{2 to 1 per unit} \\ \text{time by stimulated} \\ \text{emission} \\ \left. \vphantom{\frac{\text{\# atoms going from}}{\text{2 to 1 per unit}}} \right\} \\ + \end{array} \begin{array}{c} \text{\# atoms going from} \\ \text{2 to 1 per unit time} \\ \text{by spontaneous emission} \\ \left. \vphantom{\frac{\text{\# atoms going from}}{\text{2 to 1 per unit time}}} \right\} \end{array}$$

$$\boxed{N_1 \cdot B_{12} U(w, T) = N_2 \cdot B_{21} U(w, T) + N_2 \cdot A_{21}} \quad (36)$$

Eq.(36)  $\rightarrow \frac{N_2}{N_1} = \frac{B_{12} U(w, T)}{B_{21} U(w, T) + A_{21}} = e^{-\frac{\hbar w}{kT}}$

↑ (statistical physics) (Eq.(35))

Solve for  $U(w, T)$

$$\boxed{U(w, T) = \frac{A_{21}}{B_{12} e^{\frac{\hbar w}{kT}} - B_{21}}}$$

(37) [started to look like (33)]

Compare  $U(\omega, T) = \frac{A_{21}}{B_{12} e^{\frac{\hbar\omega}{kT}} - B_{21}}$  with  $U(\omega, T) = \frac{\hbar\omega^3}{\pi c^3} \frac{1}{e^{\frac{\hbar\omega}{kT}} - 1}$  (33)

Important Results

To get Planck's formula...

(a)  $B_{12} = B_{21} = B$  (38) [same as obtained by QM  $\lambda_{1 \rightarrow 2} = \lambda_{2 \rightarrow 1}$  (same  $U(\omega_{12})$ )]

(b)  $U(\omega, T) = \frac{A_{21}}{B} \cdot \frac{1}{e^{\frac{\hbar\omega}{kT}} - 1} \Rightarrow \boxed{\frac{A_{21}}{B} = \frac{A}{B} = \frac{\hbar\omega^3}{\pi^2 c^3}} \propto \omega^3$  (39)

$$\Rightarrow \boxed{A = \frac{\hbar\omega^3}{\pi^2 c^3} B}$$

Yes, we do!  
From QM  
(Einstein didn't)

(c) Need to introduce  $B_{21}$  (stimulated emission, not known in 1917)  
to get "-1" in Planck's formula

$$(d) A_{21} \propto \omega^3 B$$

- implies a state high in energy (so  $\omega_{21}$  is big) will "decay" to state "1" much readily (in shorter time) than a state with energy closer (thus smaller  $\omega_{21}$ ) to state "1"
- thus, harder to "hold" an atom in a highly excited state (this is true) due to faster spontaneous emission
- OR a highly excited state has a shorter life time

$$(e) A_{21} = \frac{\hbar \omega^3}{\pi^2 C^3} B \Rightarrow A_{21} \text{ & } B \text{ have different units}$$

By Eq.(34b), A has unit of  $\frac{1}{\text{time}}$  or  $\boxed{\frac{1}{A} \text{ is a time}}$

(f) [Optional]

$$A_{21} = \frac{\hbar\omega^3}{\pi^2 C^3} B = \frac{(\hbar\omega) \cdot \omega^2}{\pi^2 C^3} B$$

In statistical physics/statistical mechanics/solid state physics courses, we will see  $(\hbar\omega^3)$  actually comes from

$(\hbar\omega) \cdot \omega^2$

↑      ↗

an energy  
(quantum mechanics)

from counting number of modes  
with angular freq.  $\omega \rightarrow \omega + d\omega$   
[which is  $\sim \omega^2 d\omega$ ]